

Unification and light vectorlike fermions

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Outline

Insensitive unification of gauge couplings:

- understanding values of gauge couplings from particle content
- extensions of the SM or MSSM by complete vector-like families

Possible phenomenological consequences:

- muon $g-2$, $h \rightarrow \mu\mu$, $h \rightarrow \gamma\gamma$, $h \rightarrow ZZ^*$, $h \rightarrow WW^*$, ...
- new (possibly discovery) decay modes for heavy Higgs bosons

Proton decay

Gauge couplings in the standard model

$$\alpha_3(M_Z)_{exp} = 0.1184$$

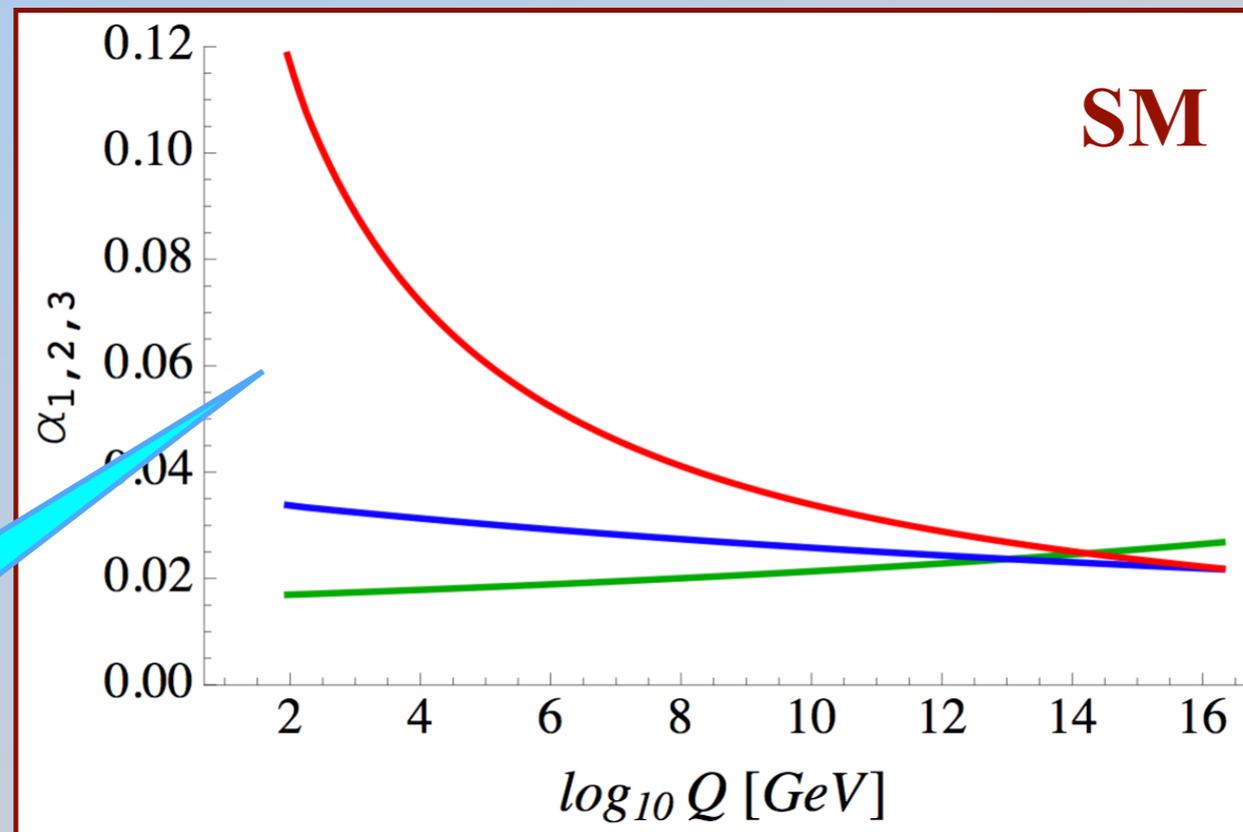
$$\alpha_2(M_Z)_{exp} = 0.03380$$

$$\alpha_1(M_Z)_{exp} = 0.01695$$

$$\alpha_{EM}(M_Z) = 1/127.916$$

$$\sin^2 \theta_W = 0.2313$$

couplings run away from each other, ratios sensitive to boundary conditions

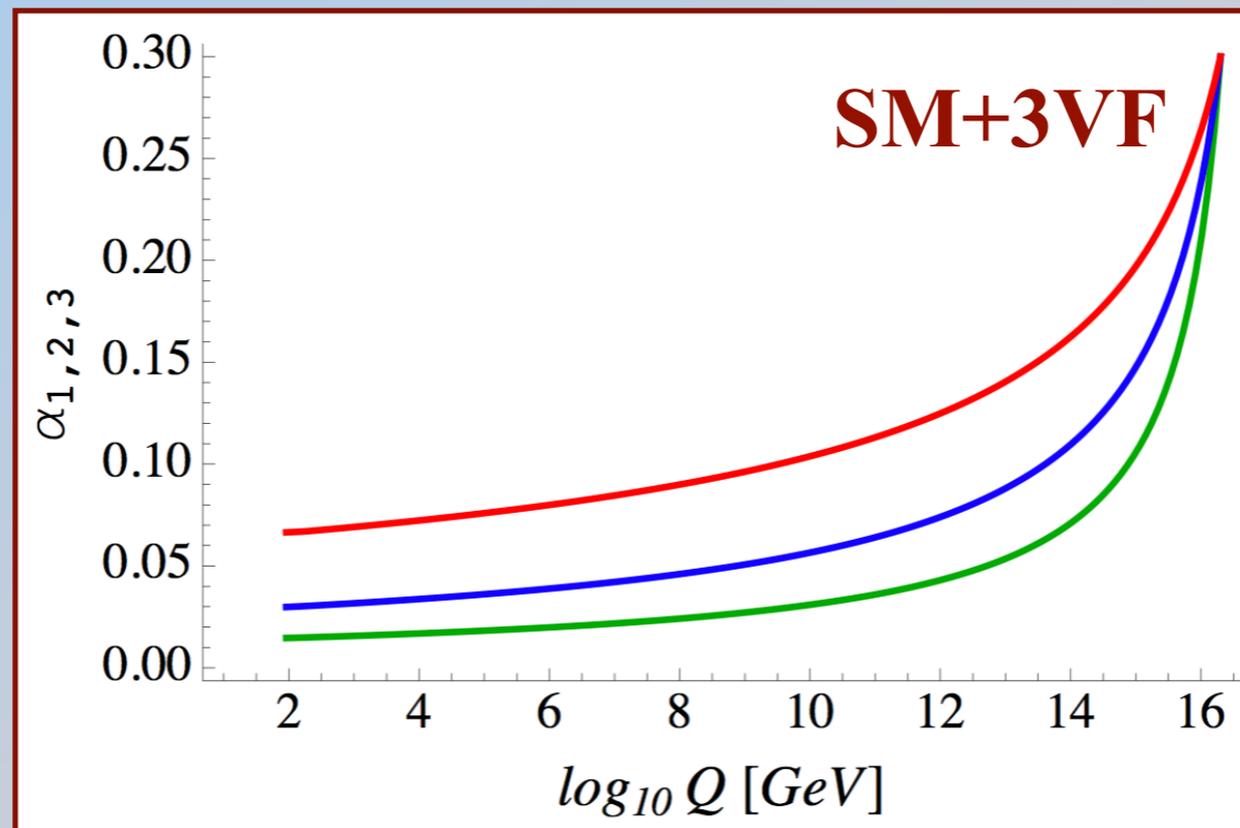


solution to RGEs:

$$\alpha_i^{-1}(M_Z) = \frac{b_i}{2\pi} \ln \frac{M_G}{M_Z} + \alpha_i^{-1}(M_G)$$

$$b_i = (41/10, -19/6, -7)$$

Gauge couplings in the SM + 3VFs



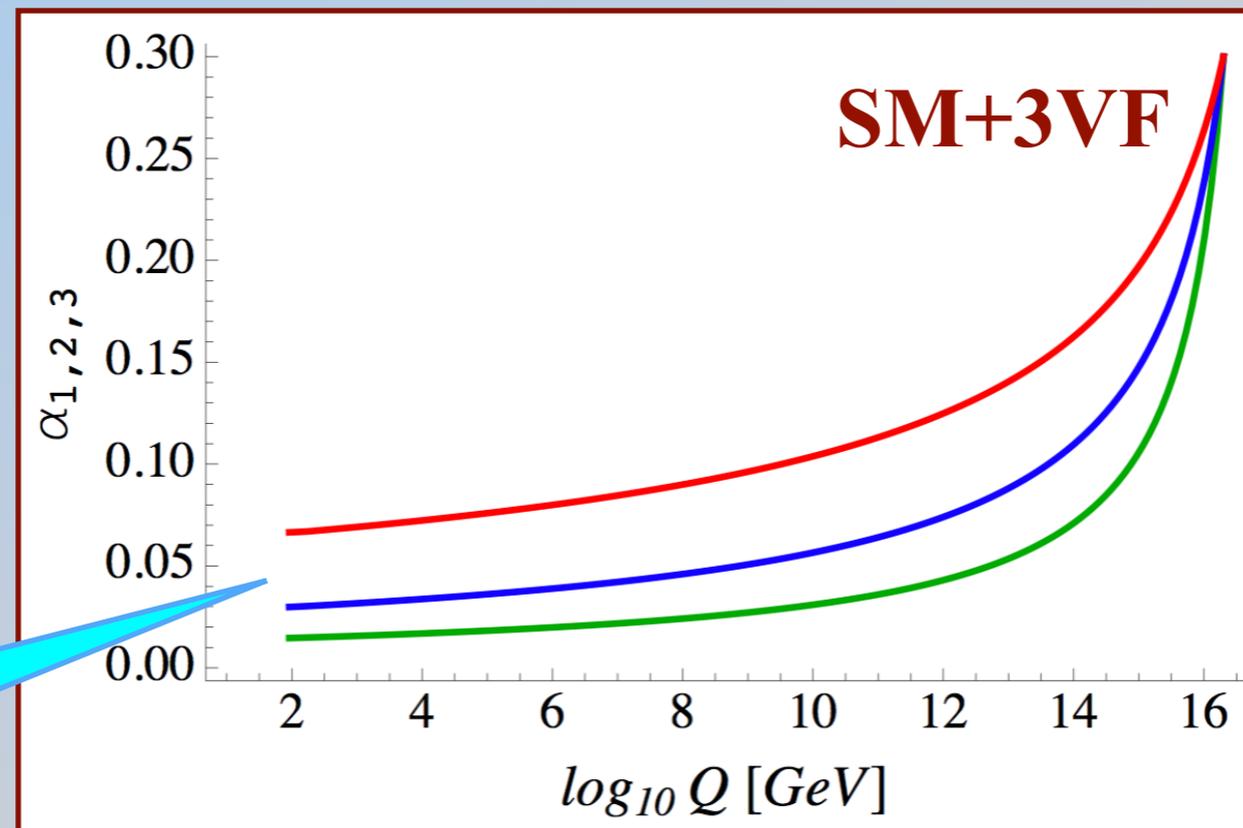
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Gauge couplings in the SM + 3VFs

couplings run
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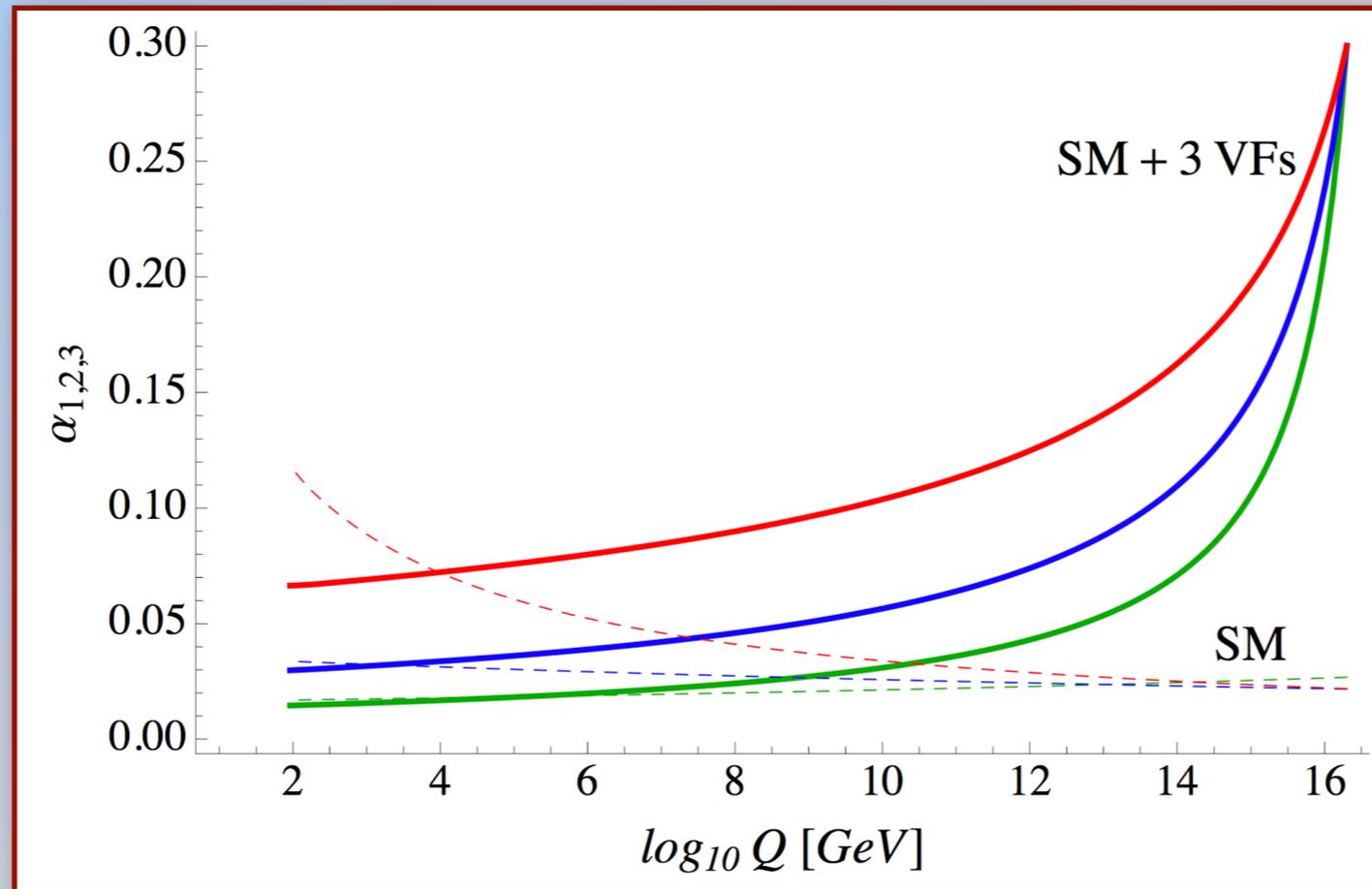


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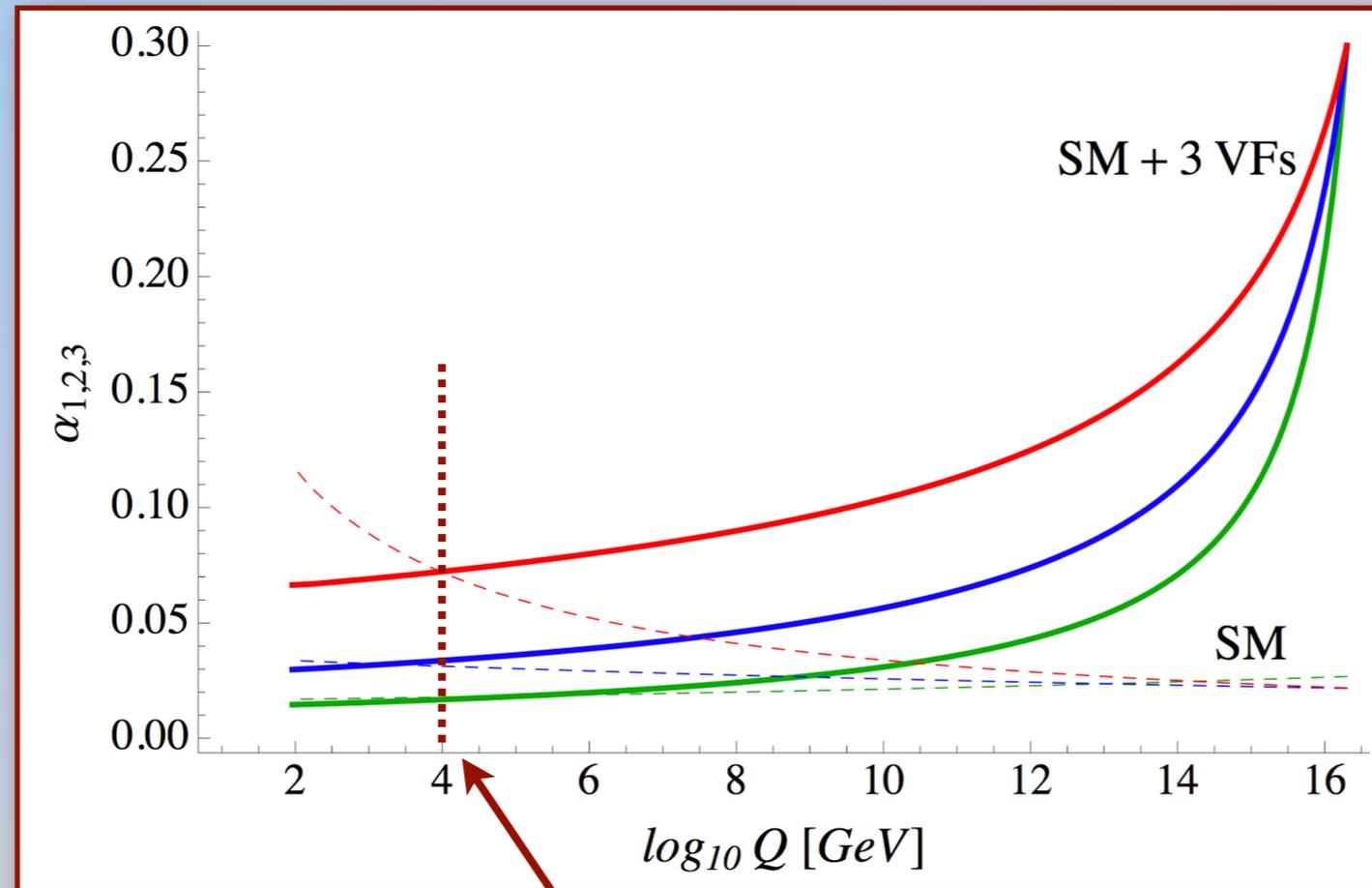
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Gauge couplings in the SM + 3 VFs

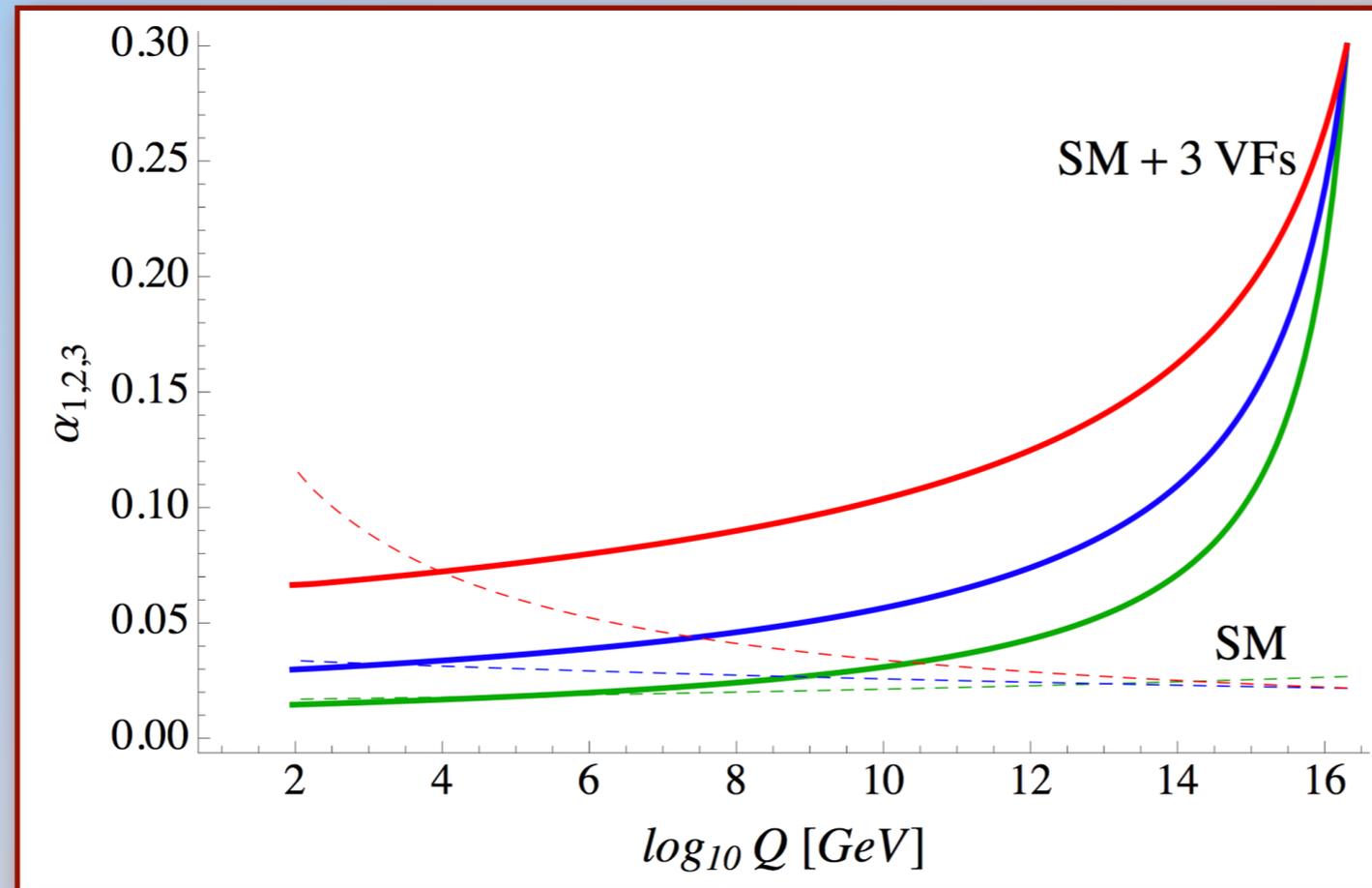


Gauge couplings in the SM + 3 VFs



Is this a threshold effect?

Gauge couplings in the SM + 3 VFs



$$\frac{\alpha_i(M_Z)}{\alpha_j(M_Z)} \simeq \frac{b_j}{b_i}$$

gauge couplings understood from:

$$\alpha_i^{-1}(M_Z) = \frac{b_i}{2\pi} \ln \frac{M_G}{M_Z} + \alpha_i^{-1}(M_G)$$

● IR fixed point predictions (two parameter free predictions)

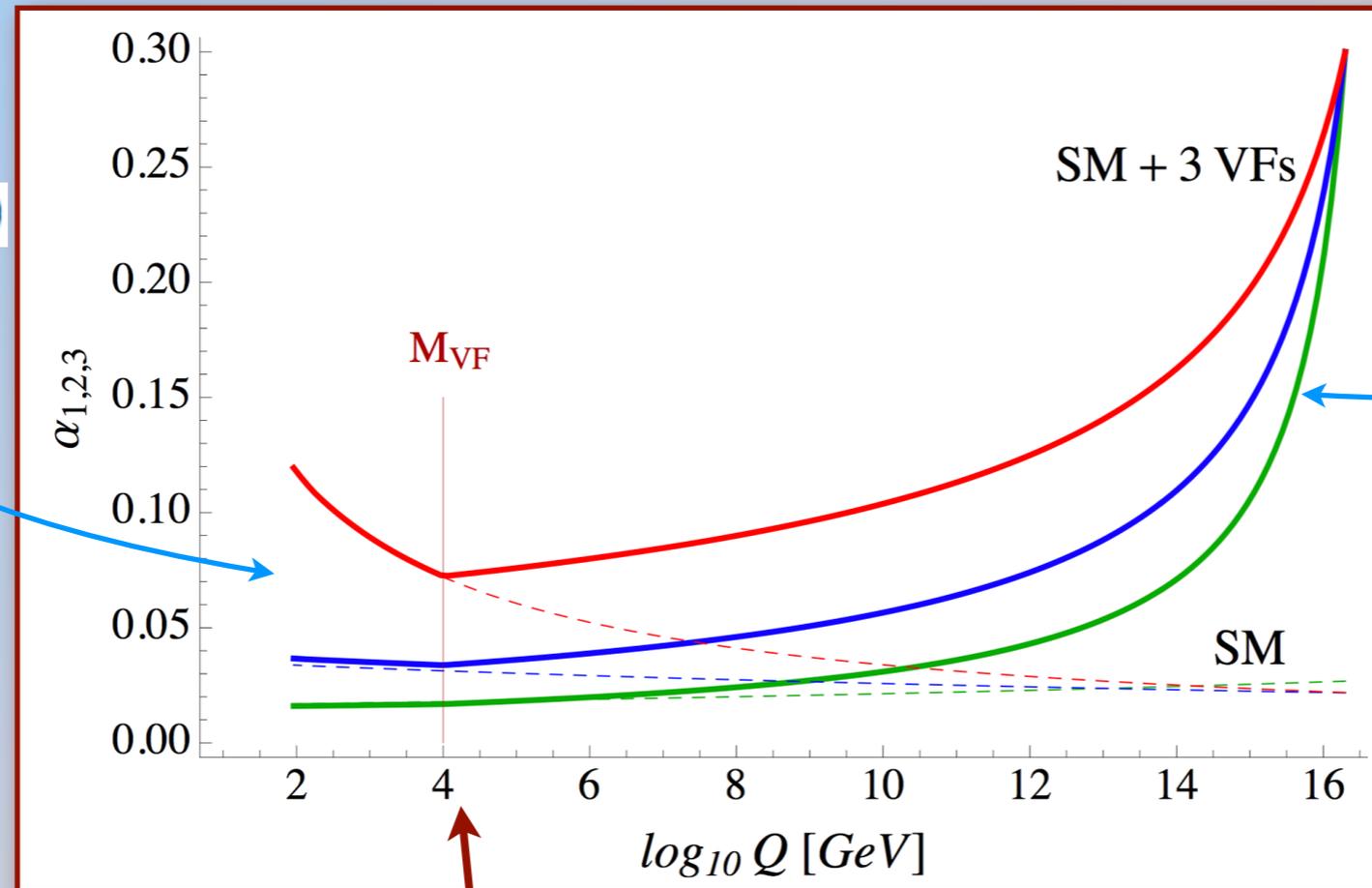
$$\sin^2 \theta_W \equiv \frac{\alpha'}{\alpha_2 + \alpha'} = \frac{b_2}{b_2 + b'} = 0.193$$

Maiani, Parisi, and Petronzio (1978)

$$\alpha_3 |_{\alpha_{EM}^{exp}} \simeq 0.072$$

(includes 2-loop)

Gauge couplings in the SM + 3 VFs



$$b_i = (41/10, -19/6, -7)$$

$$b_i = (121/10, 29/6, +1)$$

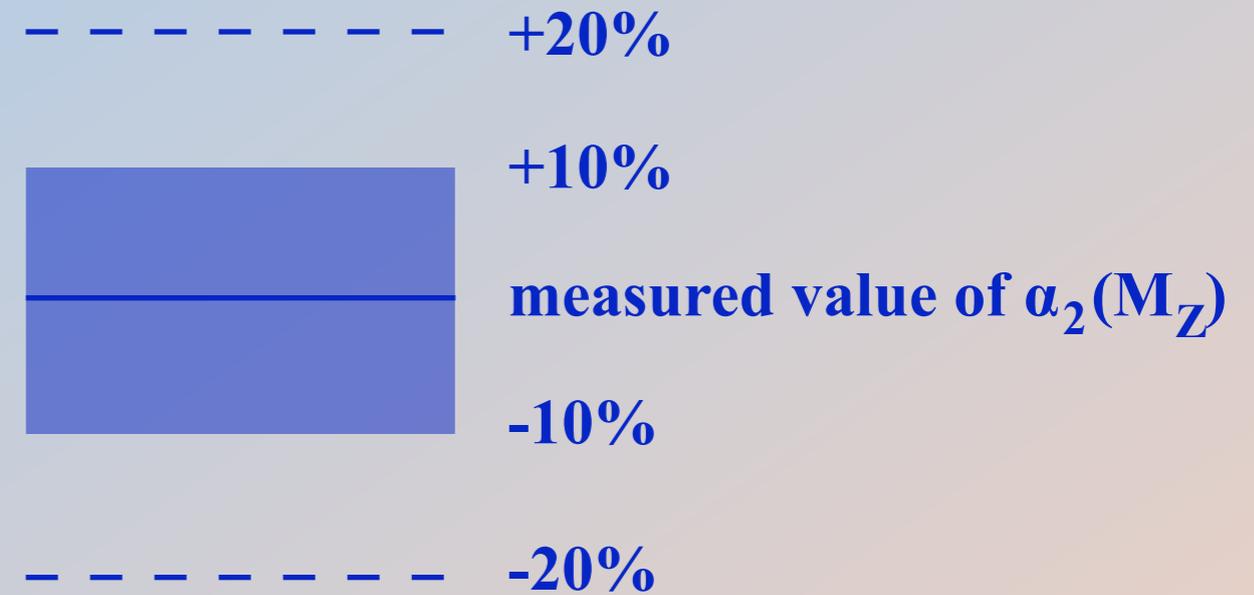
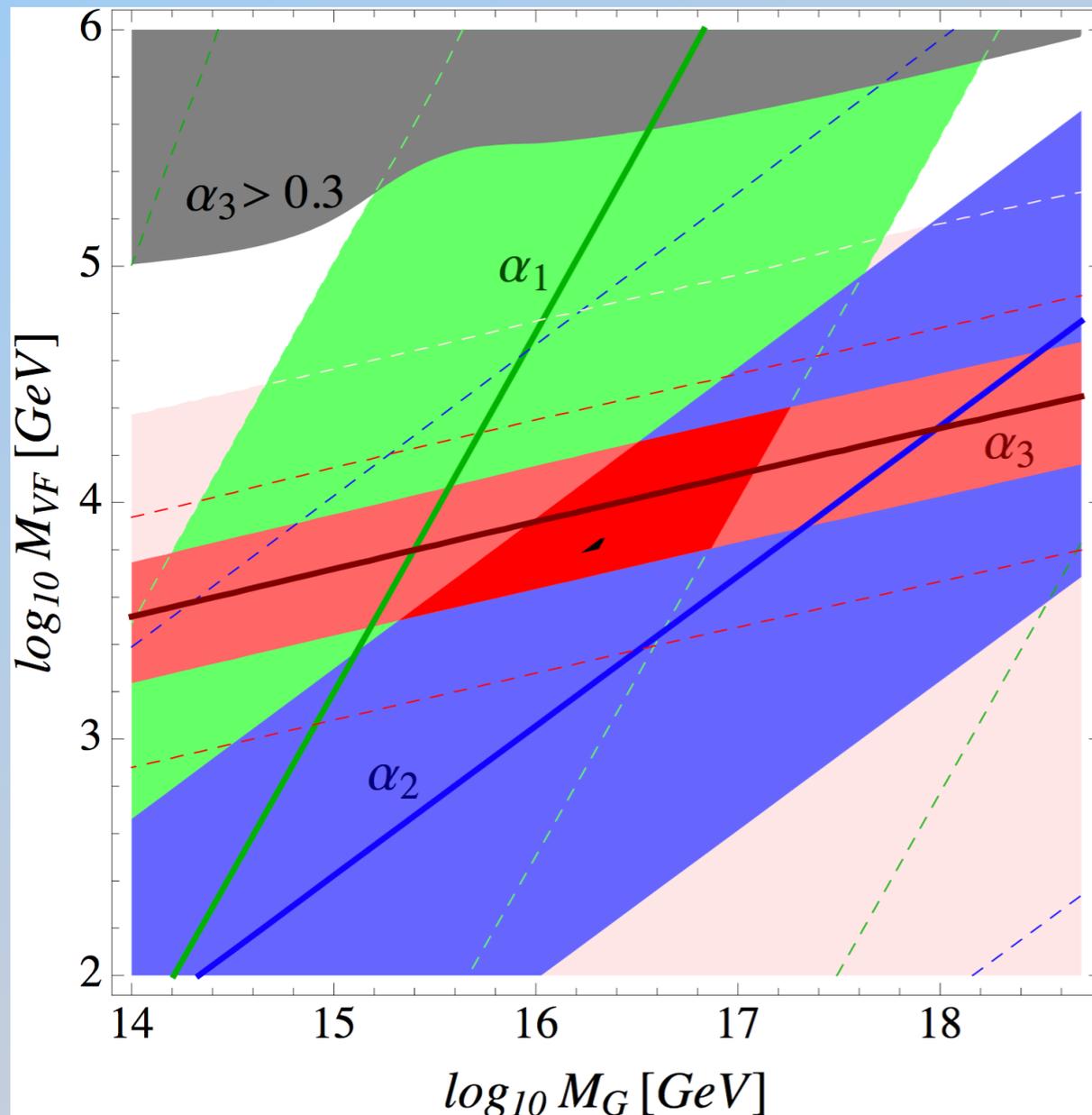
gauge couplings understood from:

- IR fixed point predictions (two parameter free predictions)
- threshold effects from masses of VFs

$$\alpha_i^{-1}(M_Z) = \frac{b_i}{2\pi} \ln \frac{M_G}{M_Z} + \alpha_i^{-1}(M_G)$$

Ranges of M_G and M_{VF}

$\alpha_G = 0.3$ (for larger values results almost identical)



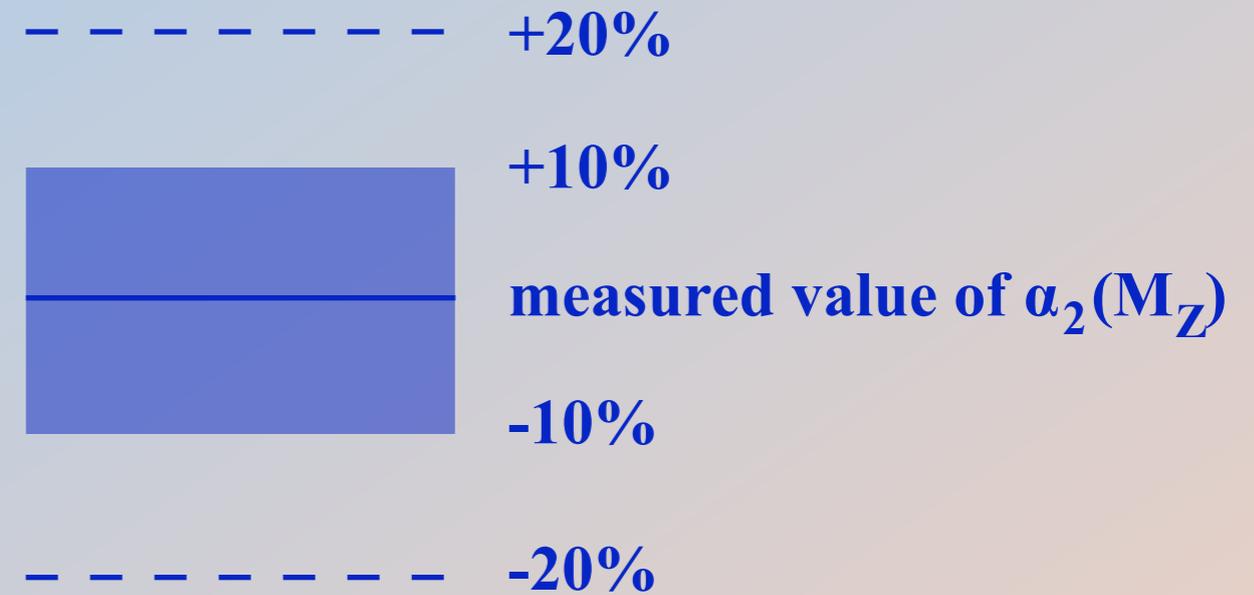
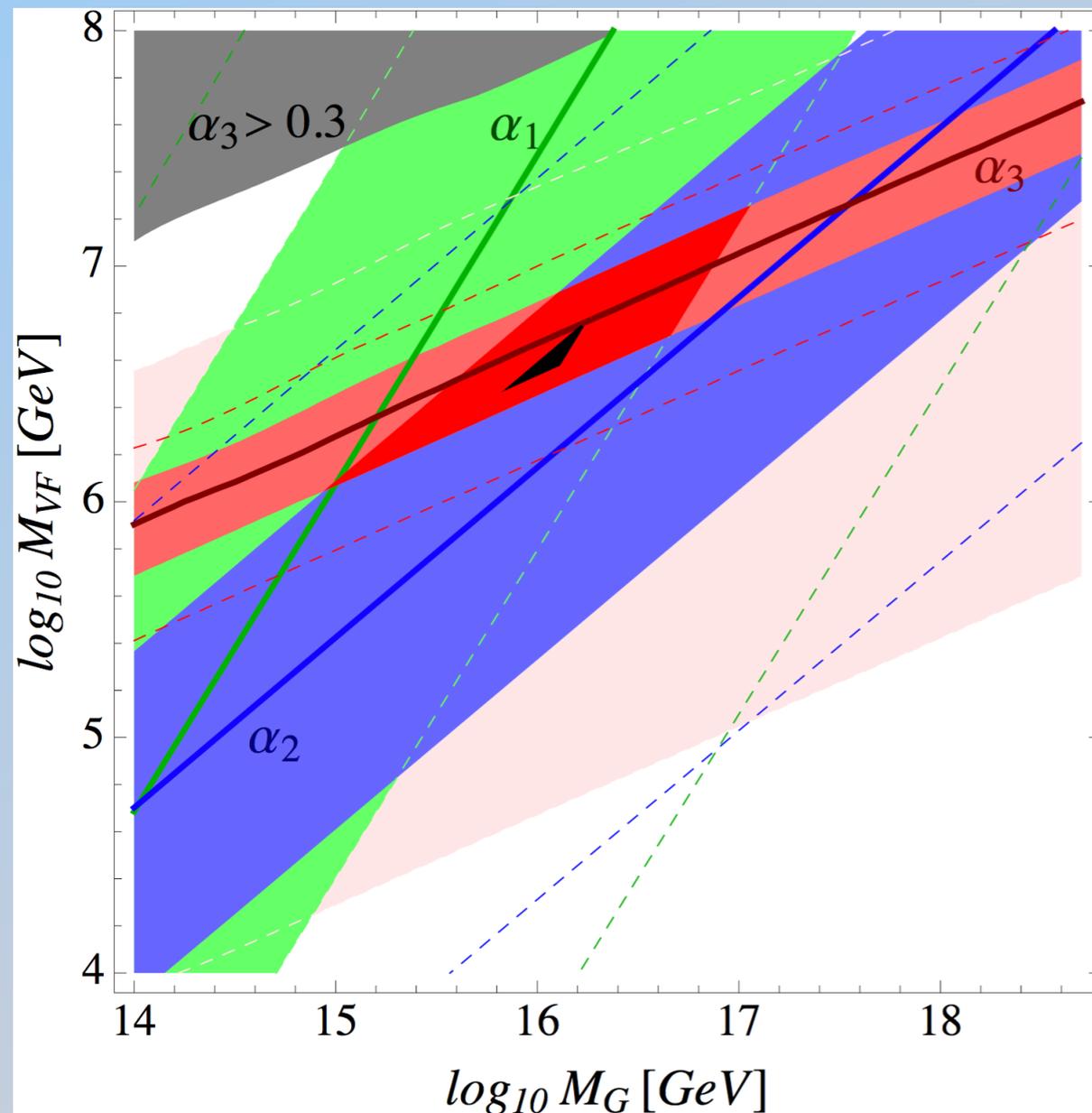
● the best fit
(all three couplings within 6%)

**Gauge couplings at M_Z within 10%
in a large range of parameters!**

**And within 50% in basically the
whole range!**

Ranges of M_G and M_{VF} for SM+4VFs

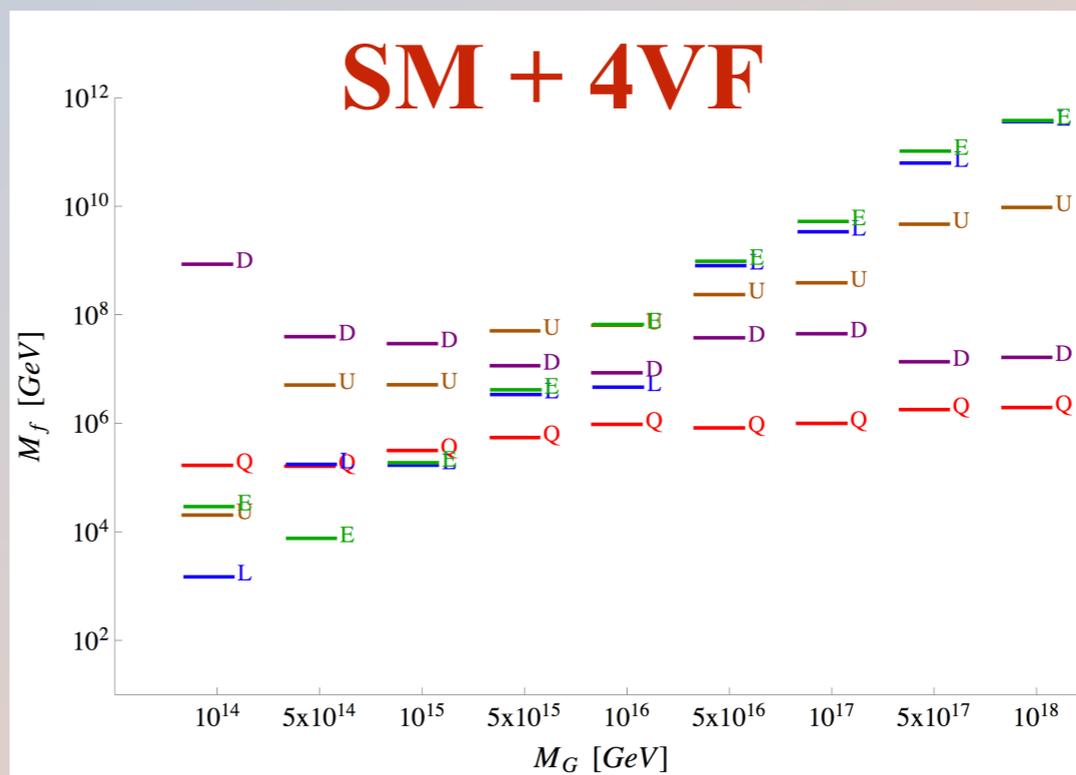
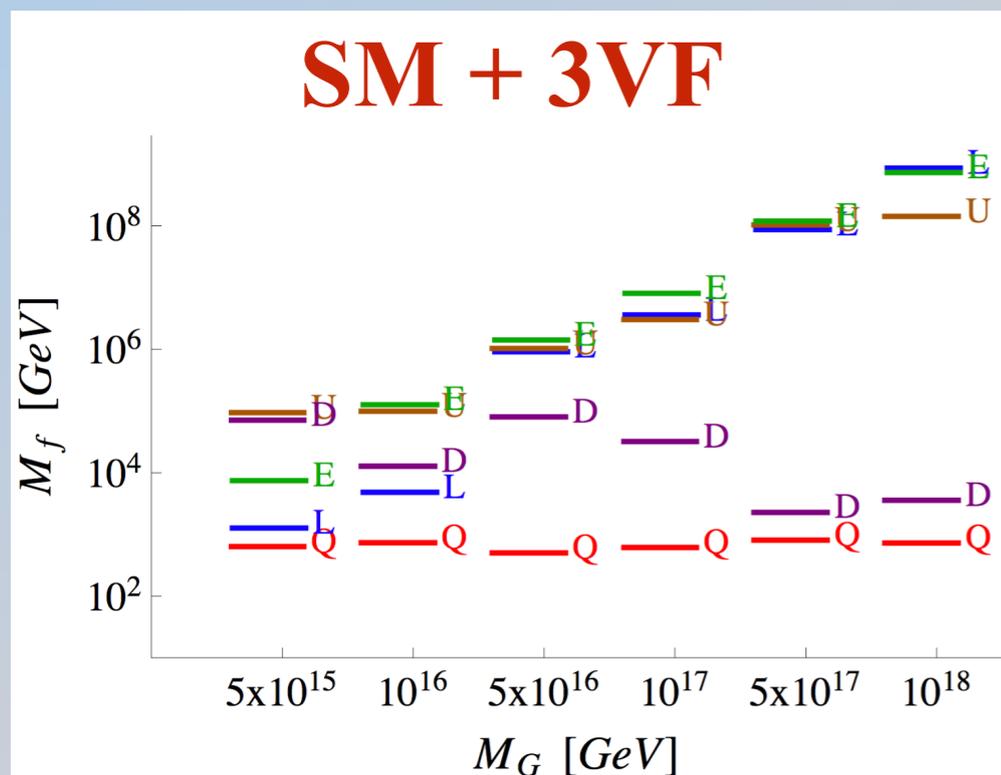
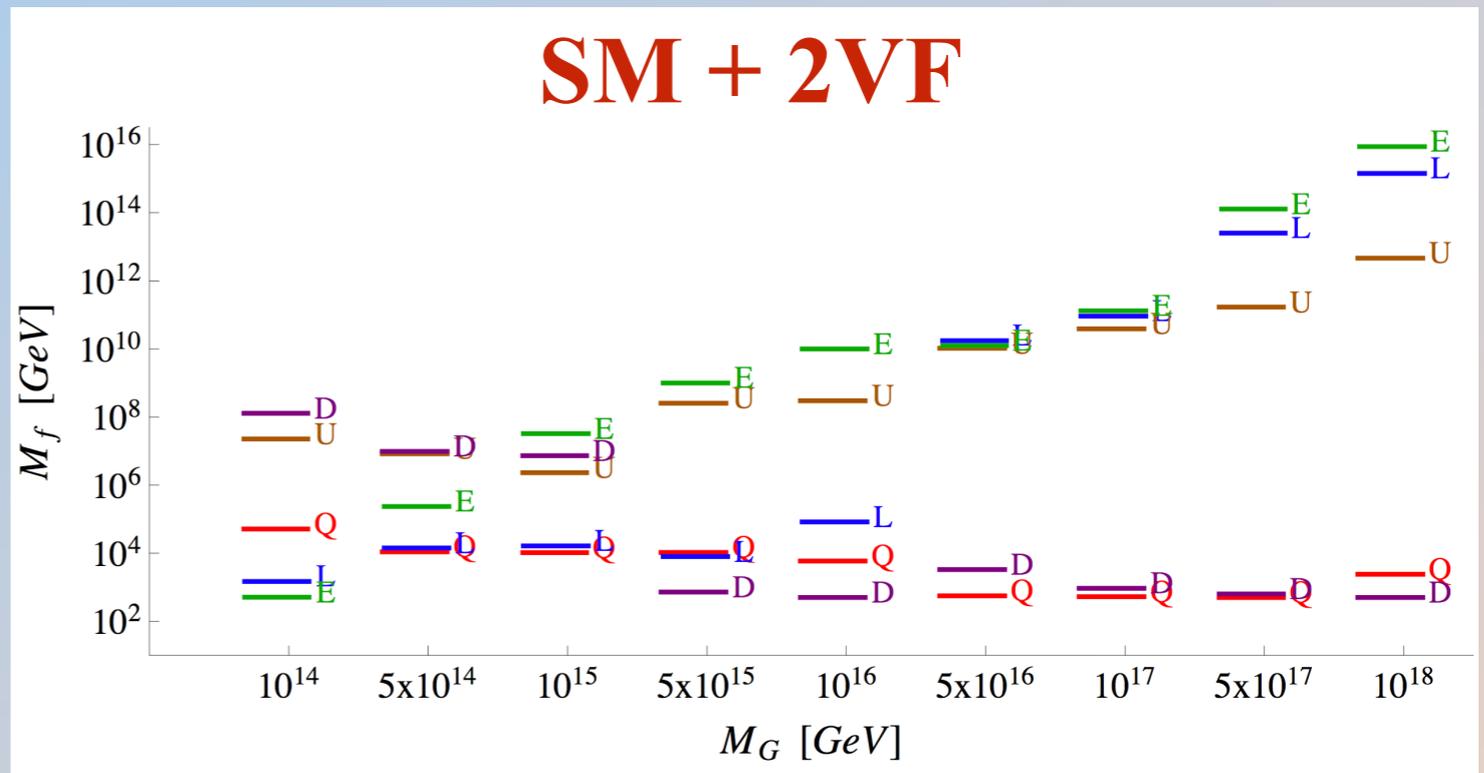
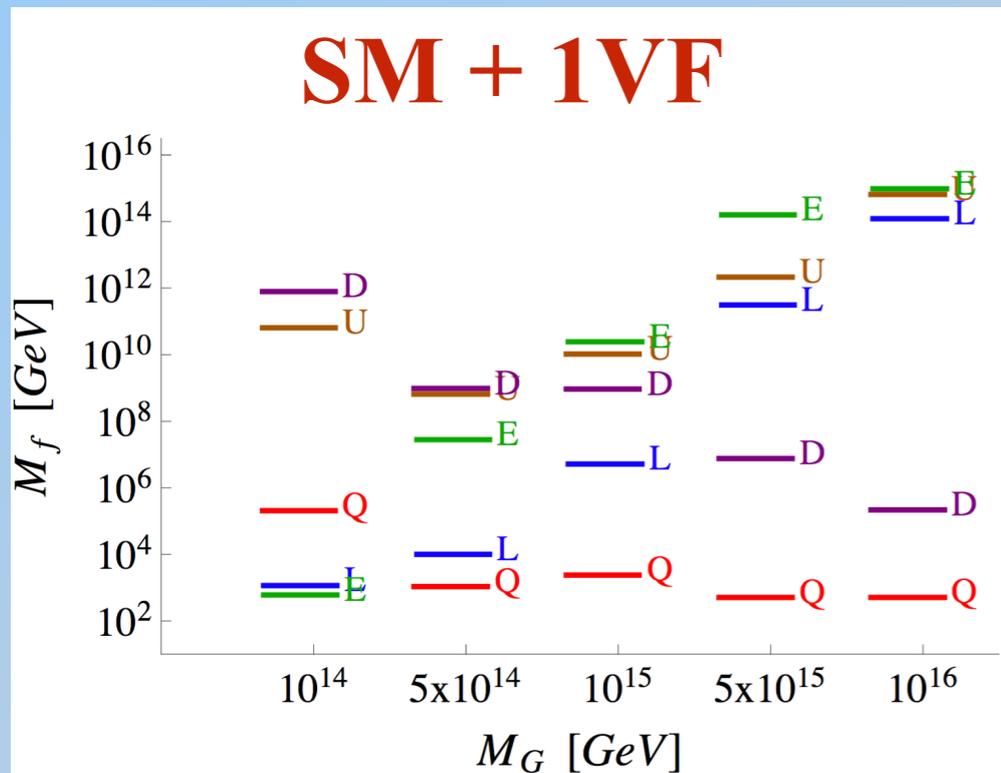
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● the best fit
(all three couplings within 6%)

**Gauge couplings at M_Z within 10%
in a large range of parameters!**

Spectrum consistent with g.c.u.



Realistic example

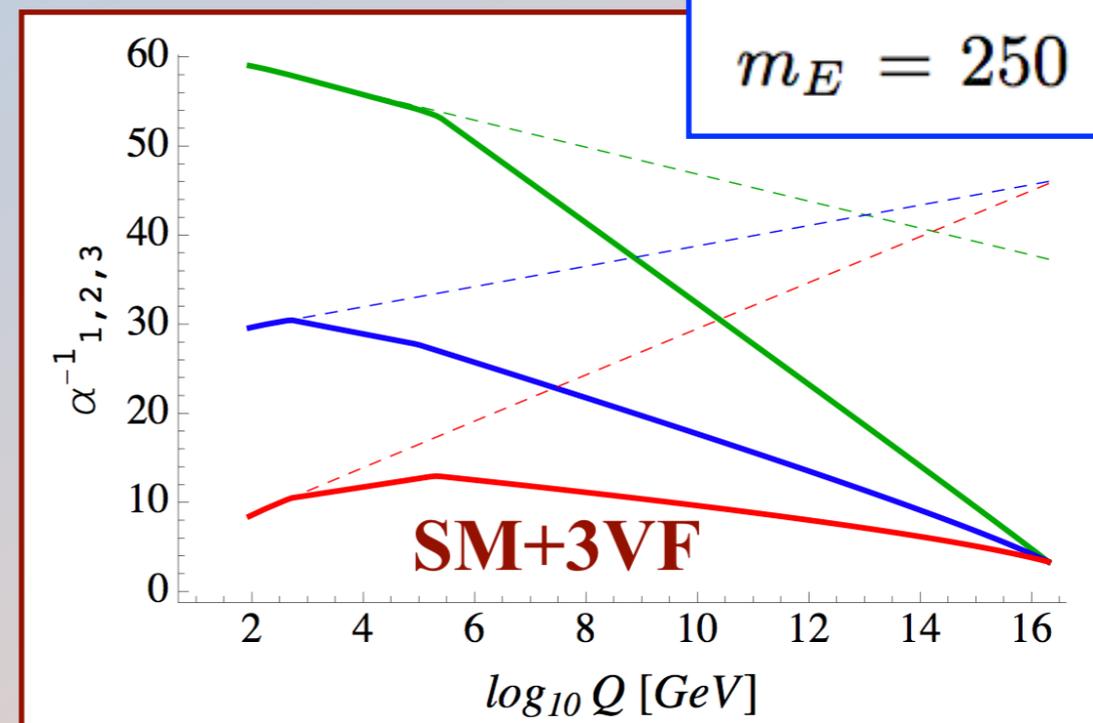
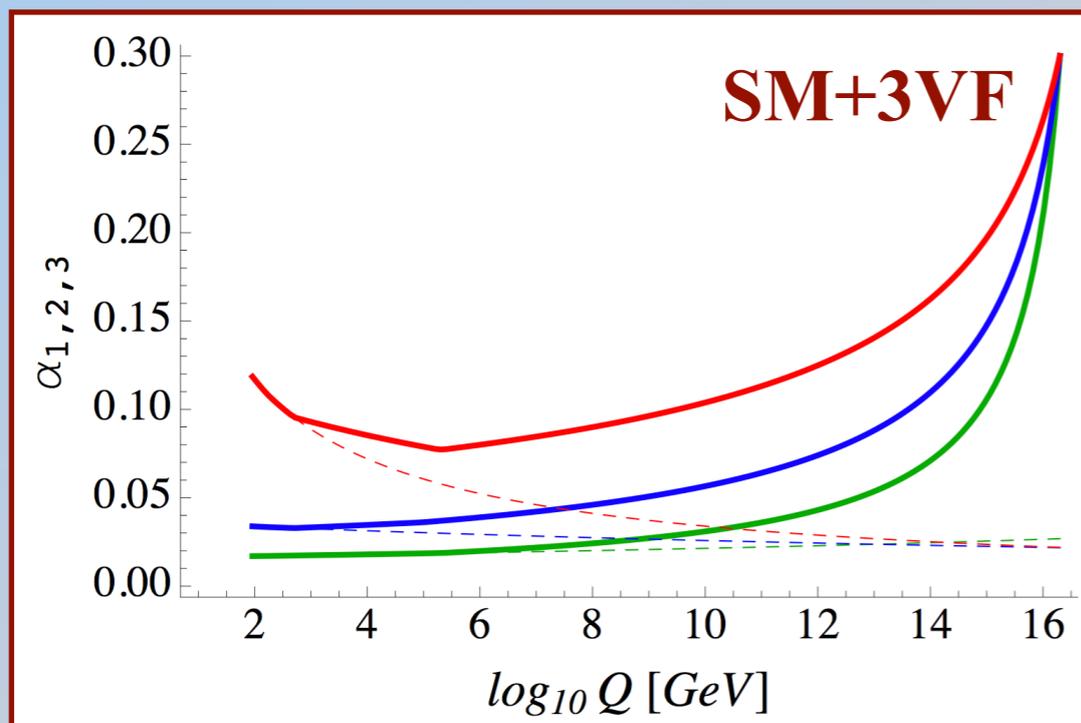
$$\begin{aligned} \alpha_3(M_Z)_{exp} &= 0.1184 \\ \alpha_2(M_Z)_{exp} &= 0.03380 \\ \alpha_1(M_Z)_{exp} &= 0.01695 \\ \alpha_{EM}(M_Z) &= 1/127.916 \\ \sin^2 \theta_W &= 0.2313 \end{aligned}$$

Gauge couplings reproduced (within fractions of exp. uncertainties) for :

4 sig. figures

2 sig. figures

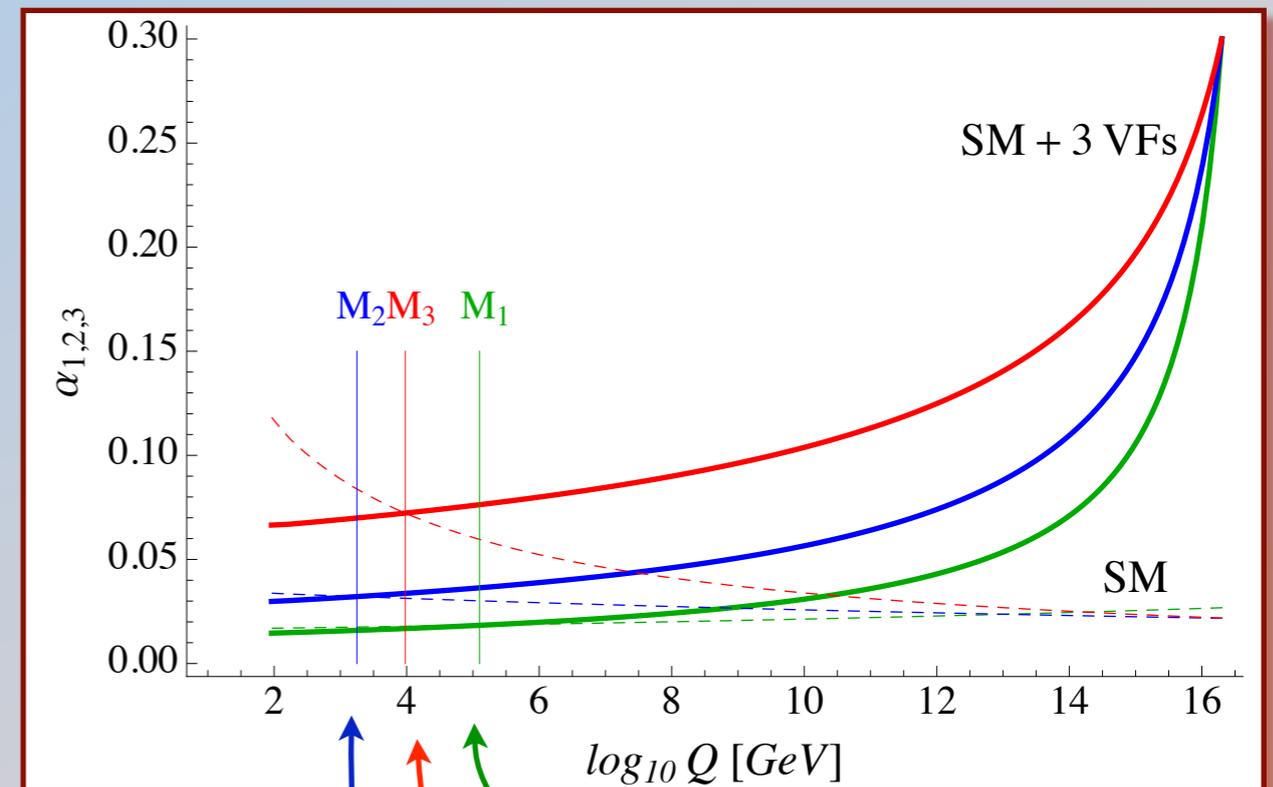
$$\begin{aligned} m_Q &= 500 \text{ GeV} \\ m_L &= 95 \text{ TeV} \\ m_U &= 220 \text{ TeV} \\ m_D &= 180 \text{ TeV} \\ m_E &= 250 \text{ TeV} \end{aligned}$$



Many possible solutions!

Classifying solutions - mass rules

To get gauge coupling unification, weighted sum of logs of masses of particles charged under given symmetry must be as if all particles had the mass equal to the crossing scale of RG evolutions of the gauge coupling in the SM and SM+3VFs:



$$\frac{1}{2\pi} \sum_{i=1}^3 \left(b_3^Q \ln \frac{M_{Qi}}{M_Z} + b_3^U \ln \frac{M_{Ui}}{M_Z} + b_3^D \ln \frac{M_{Di}}{M_Z} \right) = \frac{4}{\pi} \ln \frac{M_3}{M_Z}$$

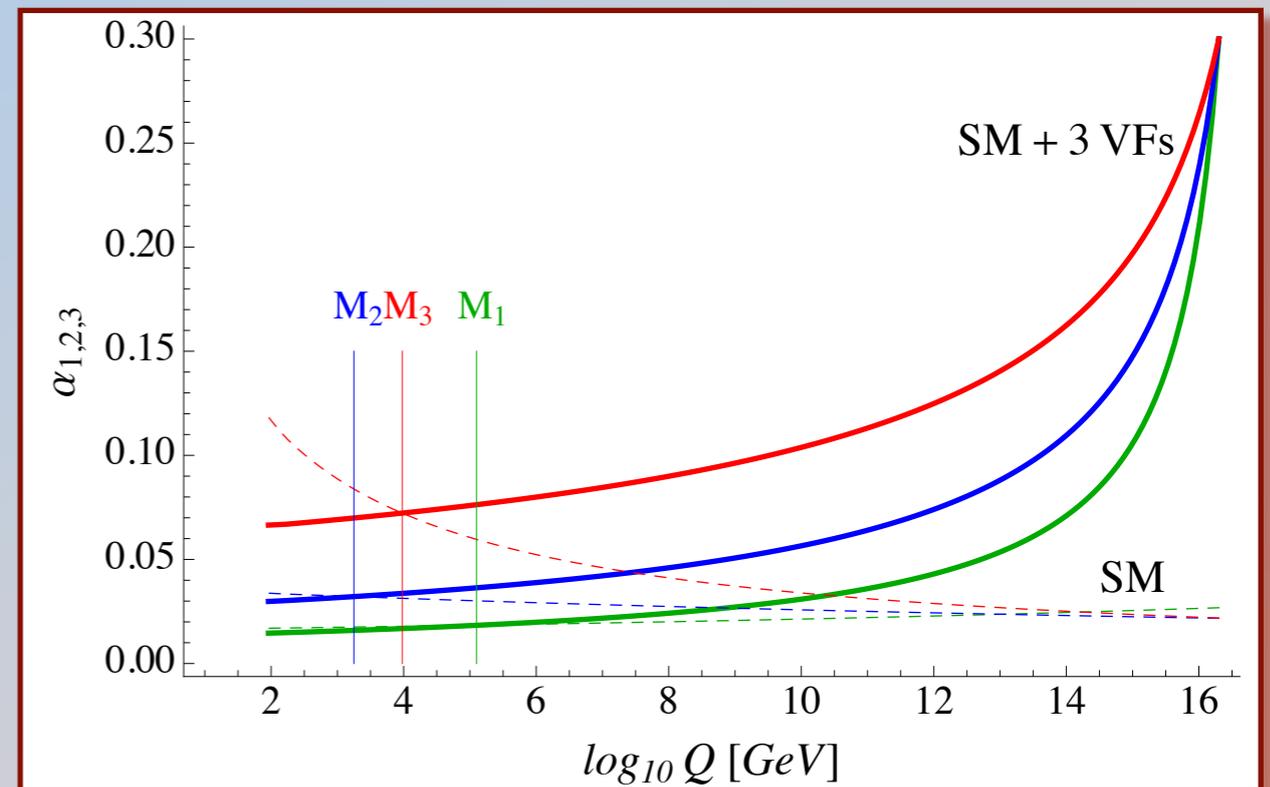
$$\frac{1}{2\pi} \sum_{i=1}^3 \left(b_2^Q \ln \frac{M_{Qi}}{M_Z} + b_2^L \ln \frac{M_{Li}}{M_Z} \right) = \frac{4}{\pi} \ln \frac{M_2}{M_Z}$$

$$\frac{1}{2\pi} \sum_{i=1}^3 \left(b_1^Q \ln \frac{M_{Qi}}{M_Z} + b_1^U \ln \frac{M_{Ui}}{M_Z} + b_1^D \ln \frac{M_{Di}}{M_Z} + b_1^L \ln \frac{M_{Li}}{M_Z} + b_1^E \ln \frac{M_{Ei}}{M_Z} \right) = \frac{4}{\pi} \ln \frac{M_1}{M_Z}$$

$$T_i = \frac{1}{2\pi} \sum_f b_i^f \ln \frac{M_f}{M_Z}$$

Classifying solutions - mass rules

To get gauge coupling unification, weighted sum of logs of masses of particles charged under given symmetry must be as if all particles had the mass equal to the crossing scale of RG evolutions of the gauge coupling in the SM and SM+3VFs:



$$M_3^4 = M_Q^2 M_U M_D,$$

$$M_2^4 = M_Q^3 M_L,$$

$$M_1^{20} = M_Q M_U^8 M_D^2 M_L^3 M_E^6.$$

$$M_F \equiv (M_{F_1} M_{F_2} \dots M_{F_N})^{1/N}$$

What else are extra VFs good for?

◆ muon $g-2$

K. Kannike, M. Raidal, D.M. Straub and A. Strumia, 1111.2551 [hep-ph]

R.D. and A. Raval, 1305.3522 [hep-ph]

$$\Delta a_\mu = a_\mu^{exp} - a_\mu^{SM} = 2.7 \pm 0.80 \times 10^{-9}$$

◆ anomalies in Z-pole observables A_{FB}^b and A_e

D. Choudhury, T.M.P. Tait and C.E.M. Wagner, hep-ph/0109097

R.D., S.G. Kim and A. Raval, 1105.0773 [hep-ph], 1201.0315 [hep-ph]

B. Batell, S. Gori and L.T. Wang 1209.6382 [hep-ph]

◆ Higgs boson mass and EWSB

K.S. Babu, I. Gogoladze, M.U. Rehman, Q. Shafi, 0807.3055 [hep-ph]

S.P. Martin, 0910.2732 [hep-ph]

◆ $h \rightarrow \gamma\gamma, ZZ^*, WW^*, \dots$

really many papers

◆ ...

Two Higgs doublet model - type II

VL mixing only with 2nd generation of leptons:

R.D., E. Lunghi and S. Shin, 1509.04292

$$\begin{aligned} \mathcal{L} \supset & -y_\mu \bar{\mu}_L \mu_R H_d - \lambda_E \bar{\mu}_L E_R H_d - \lambda_L \bar{L}_L \mu_R H_d - \lambda \bar{L}_L E_R H_d - \bar{\lambda} H_d^\dagger \bar{E}_L L_R \\ & - \kappa_N \bar{\mu}_L N_R H_u - \kappa \bar{L}_L N_R H_u - \bar{\kappa} H_u^\dagger \bar{N}_L L_R \\ & - M_L \bar{L}_L L_R - M_E \bar{E}_L E_R - M_N \bar{N}_L N_R + \text{h.c.} , \end{aligned}$$

$$\mu_L = \begin{pmatrix} \nu_\mu \\ \mu_L^- \end{pmatrix}, \quad L_{L,R} = \begin{pmatrix} L_{L,R}^0 \\ L_{L,R}^- \end{pmatrix}, \quad H_d = \begin{pmatrix} H_d^+ \\ H_d^0 \end{pmatrix}, \quad H_u = \begin{pmatrix} H_u^0 \\ H_u^- \end{pmatrix}$$

couplings to gauge bosons are modified because SU(2) doublets mix with SU(2) singlets and couplings to Higgs are modified because of explicit vectorlike mass terms:

$$(\bar{\mu}_L, \bar{L}_L^-, \bar{E}_L) \begin{pmatrix} y_\mu v_d & 0 & \lambda^E v_d \\ \lambda^L v_d M_L & \lambda v_d & \\ 0 & \bar{\lambda} v_d M_E & \end{pmatrix} \begin{pmatrix} \mu_R \\ L_R^- \\ E_R \end{pmatrix}$$

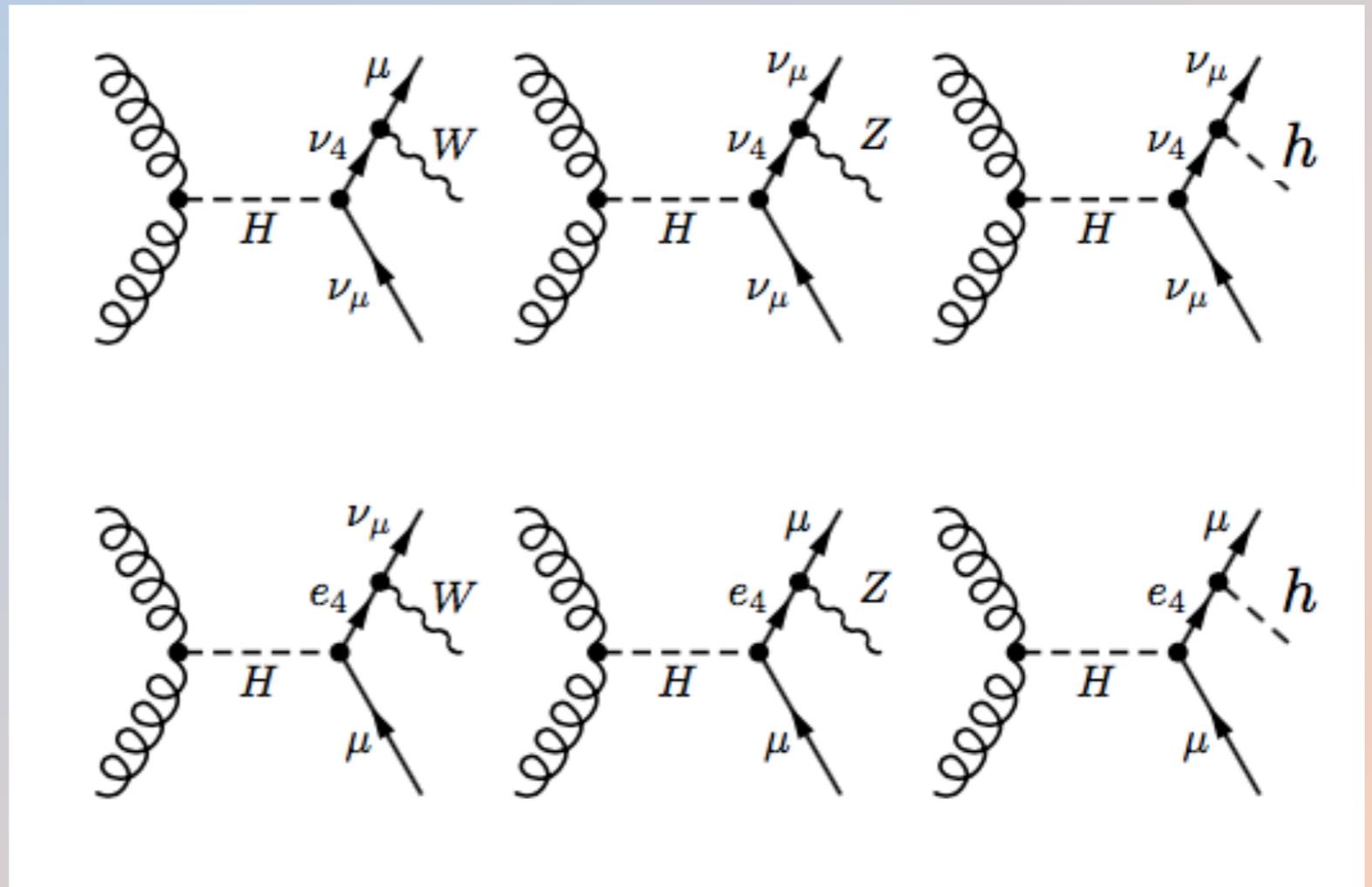
$$\begin{pmatrix} \bar{\nu}_\mu & \bar{L}_L^0 & \bar{N}_L \end{pmatrix} \begin{pmatrix} 0 & 0 & \kappa_N v_u \\ 0 & M_L & \kappa v_u \\ 0 & \bar{\kappa} v_u & M_N \end{pmatrix} \begin{pmatrix} \nu_R = 0 \\ L_R^0 \\ N_R \end{pmatrix}$$

New (possibly discovery) decay modes

$$H \rightarrow W\cancel{W}, Z\cancel{Z}, \gamma\gamma, \cancel{\tau\tau}, b\bar{b}, \tau\bar{\tau}, \dots$$

if **h** is SM-like and **H** (or **A**) is below ~ 350 GeV flavor violating decays can be dominant:

decays to pairs of heavy leptons also possible but limited to smaller mass ranges (constraints from pair-production can be applied)

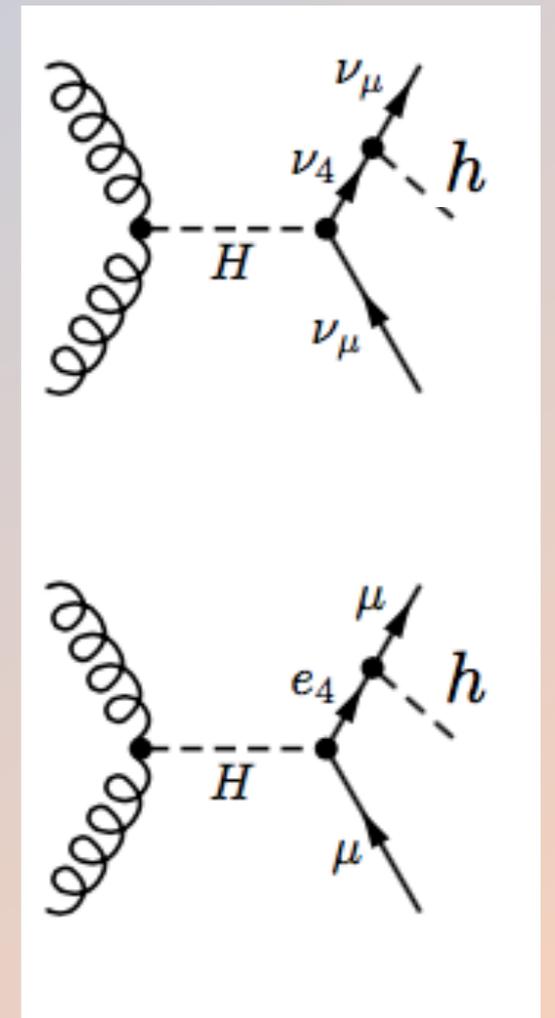


they all look similar to WW, ZZ, hZ decay modes of H or ZZ, WW, Zh production!

$H \rightarrow h \nu \nu$ and $H \rightarrow h \mu \mu$

look like Zh production, with potentially much larger cross section, (no Z , but no penalty for 2 leptons)

R.D., E. Lunghi and S. Shin, in progress

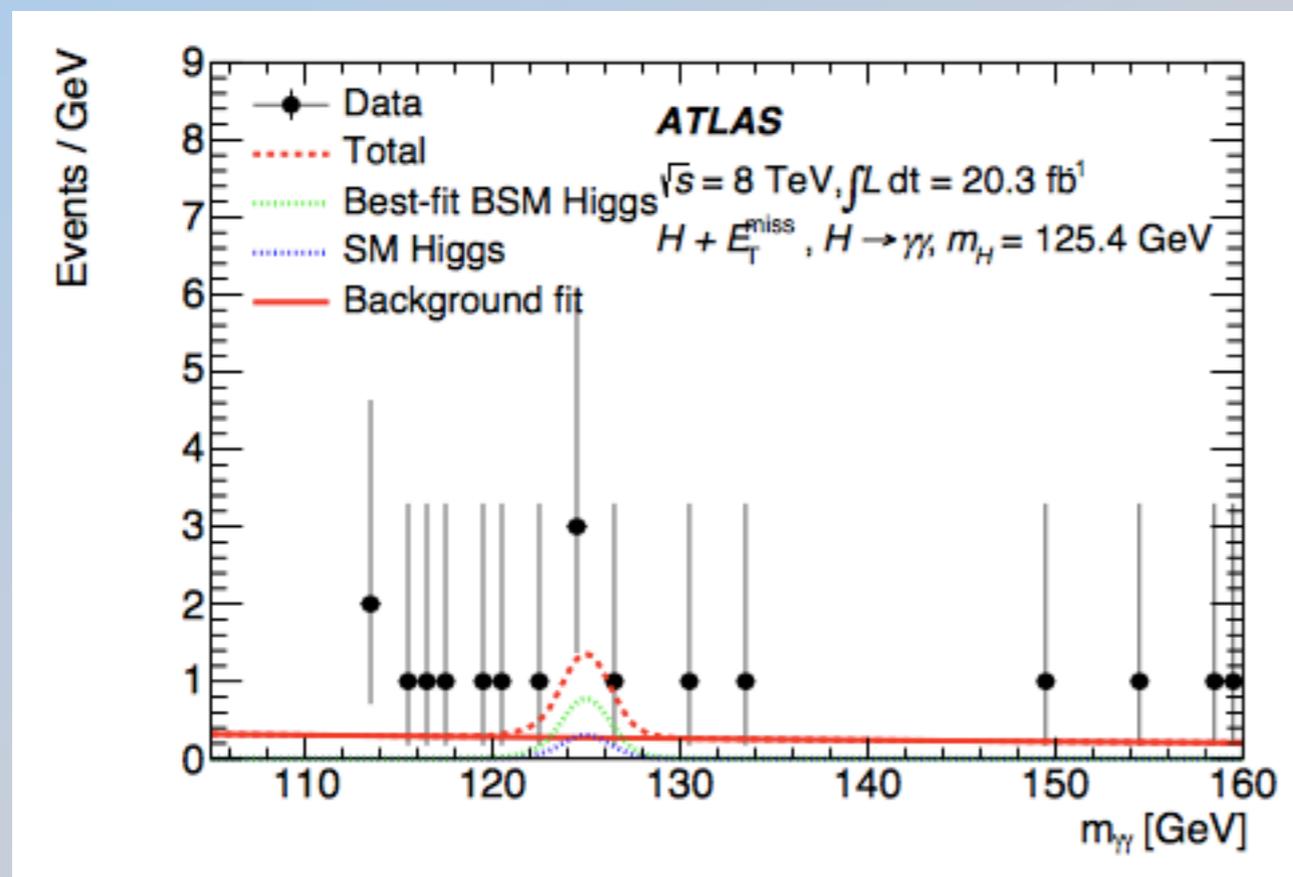


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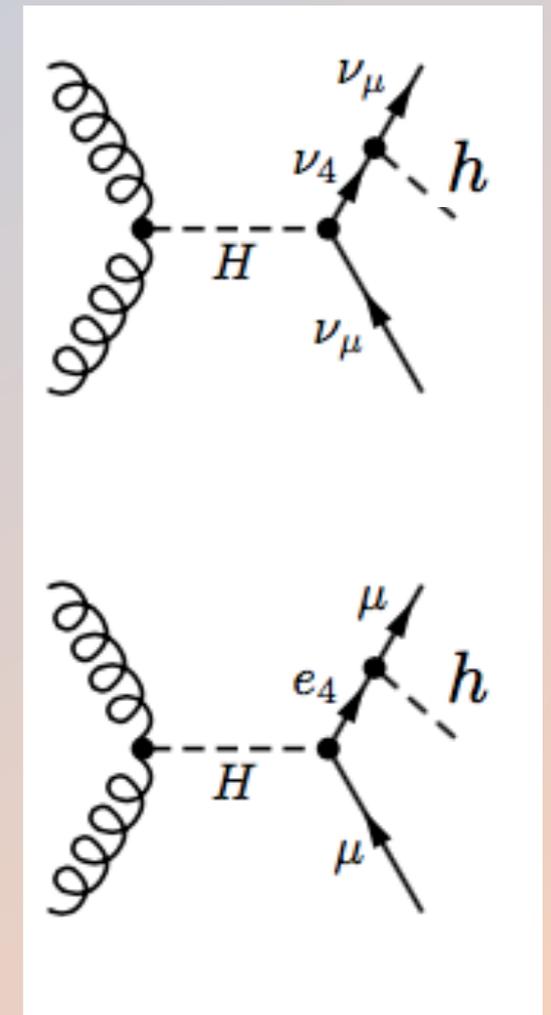
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R.D., E. Lunghi and S. Shin, in progress

some decay modes almost background free, e.g. $h \rightarrow \gamma\gamma$:

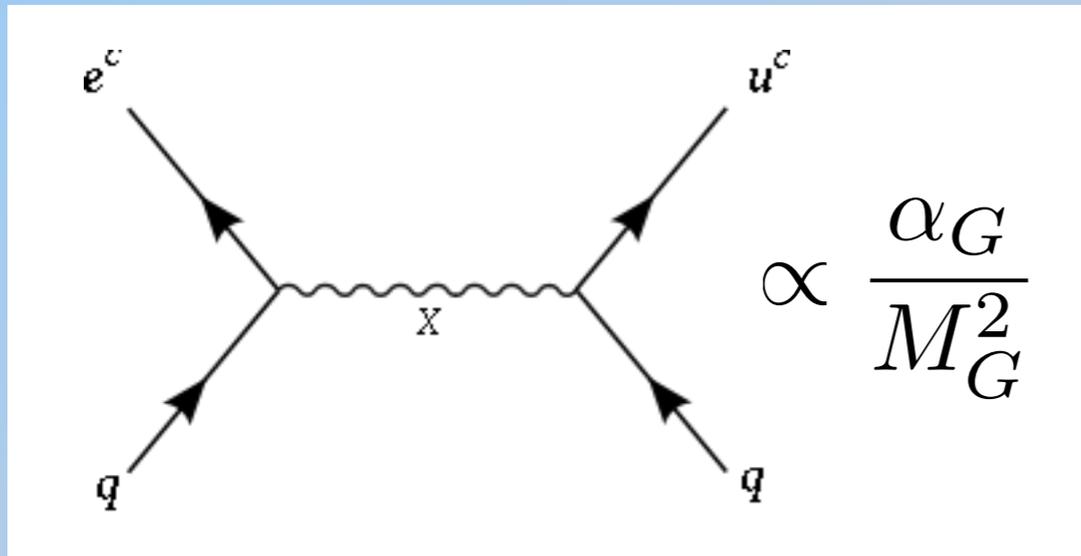


ATLAS, 1506.01081



Effects of VFs on proton lifetime

dimension 6 operators:



dominant decay mode:

$$p \rightarrow \pi^0 e^+$$

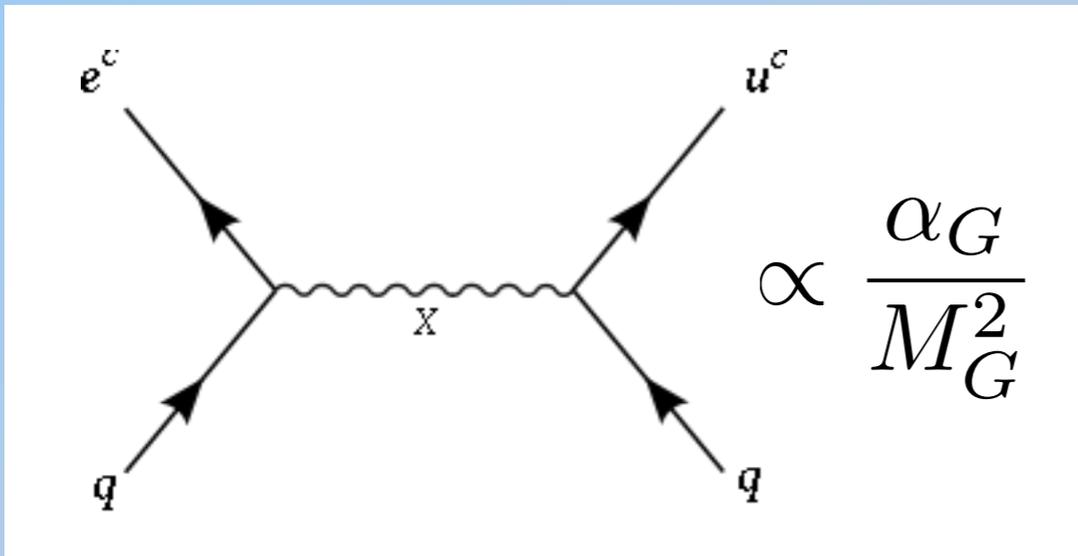
experimental limit:

$$\tau(p \rightarrow \pi^0 e^+) \gtrsim 10^{34} \text{ yrs}$$

typical prediction of SUSY GUTs: $\sim 10^{36}$ yrs

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With complete vectorlike families: the value of unified gauge coupling is significantly larger (and adjustable)

with or without SUSY the optimal GUT scale is similar

Dim. 6 operators can be close to experimental limits!

Conclusions

- ◆ **SM+3VFs, MSSM+1VF (or more) allow for insensitive unification of gauge couplings**
predictive (in a different way)
- ◆ **even without supersymmetry the optimal GUT scale is high enough to survive limits on proton lifetime**
- ◆ **the value of unified gauge coupling is larger (and adjustable)**
dim. 6 proton decay operators can be close to current limits
- ◆ **minimal scenarios predict VFs near the EW scale**
and can lead to very interesting and rich phenomenology